# **Intrabeam Scattering in the VLHC**

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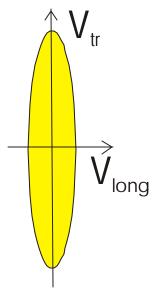
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## **VLHC Parameters**

		Stage 1	Stage 2
Beam energy	$E_p$	20 TeV	87.5 TeV
Luminosity	L	$10^{34}  \text{cm}^{-2} \text{s}^{-1}$	$2 \cdot 10^{34} \text{ cm}^{-2} \text{s}^{-1}$
Magnetic field	$B_0$	1.96 T	10 T
Injection energy	$ E_{inj} $	0.9 TeV	10 TeV
Bunch spacing	$  t_b  $	18.9 ns	
Circumference	<i>C</i> =2 <i>pR</i>	232 km	
Revolution frequency	$f_0$	1294 Hz	
Number of bunches	$N_b$	36943	
RF frequency	$f_{RF}$	477.9 MHz	
Betatron tunes	$n_{\perp}$	~214	
Momentum compaction	а	$2.1 \cdot 10^{-5}$	
Beta-function at IP	$\boldsymbol{b}^*$	30 cm	50 cm
Head-on beam-beam tune shift per IP	X	1.8·10 <sup>-2</sup>	1.8·10 <sup>-2</sup>
Number of particles per bunch	N	$2.5 \cdot 10^{10}$	
Beam current	$I_b$	0.190 A	0.175 A
RF voltage per turn, top/injection energy	$V_0$	50/50 MeV	50/50 MeV
Synchrotron frequency, top/injection energy	$ f_s $	2.32/10.9 Hz	1.1/3.28 Hz
Rms momentum spread, top/injection energy	$\sigma_{\rm p}$	1.5/14.9·10 <sup>-4</sup>	$0.5/2.4 \cdot 10^{-4}$
Rms bunch length, top/injection energy	$\sigma_{\rm s}$	6.6/14.2 cm	4.5/7.8 cm

## 1. Intrabeam scattering theory

If in the beam frame the longitudinal momentum spread is much less than the transverse one



IBS calculations can be significantly simplified

#### The IBS growth rate for longitudinal degree of freedom

• IBS transfers the energy from the transverse degrees of freedom to the longitudinal one and the growth rate can be approximated by the following formula

$$\frac{d}{dt}(\boldsymbol{q}_{\parallel}^{2}) \equiv \frac{d}{dt}\left(\frac{\overline{p}_{\parallel}^{2}}{p}\right) = \frac{e^{4}N_{i}}{4\sqrt{2}m_{p}^{2}c^{3}\boldsymbol{g}_{i}^{3}\boldsymbol{b}_{i}^{3}\boldsymbol{s}_{si}} \left\langle\frac{\Xi_{\parallel}(\boldsymbol{q}_{x},\boldsymbol{q}_{y})}{\sqrt{\boldsymbol{q}_{x}^{2}+\boldsymbol{q}_{y}^{2}}}\frac{L_{c}}{\boldsymbol{s}_{x}\boldsymbol{s}_{y}}\right\rangle_{s},$$

where averaging is performed along the beam orbit,

$$\Xi_{\parallel}(x,y) \approx 1 + \frac{\sqrt{2}}{\boldsymbol{p}} \ln \left( \frac{x^2 + y^2}{2xy} \right) - 0.055 \left( \frac{x^2 - y^2}{x^2 + y^2} \right)^2,$$

$$\boldsymbol{s}_x = \sqrt{\boldsymbol{e}_x \boldsymbol{b}_y + D^2 \boldsymbol{q}_{\parallel}^2}, \, \boldsymbol{s}_y = \sqrt{\boldsymbol{e}_y \boldsymbol{b}_y}, \, \boldsymbol{q}_x = \sqrt{\boldsymbol{e}_x / \boldsymbol{b}_x} \text{ and } \boldsymbol{q}_y = \sqrt{\boldsymbol{e}_y / \boldsymbol{b}_y}$$

- are the beam sizes and angular spreads along the ring, and
- $L_c$  is the Coulomb logarithm.
- In the smooth focusing approximation,
- for equal horizontal and vertical emittances,  $\mathbf{e}_x = \mathbf{e}_y$ ,
- equal betatron tunes,  $\mathbf{n}_x = \mathbf{n}_y$ , and
- small contribution of energy spread into the beam size,  $D(\Delta p/p) \ll \mathbf{s}_x$ ,
  - ⇒ The momentum spread growth rate can be written in the following form

$$\Lambda_{\parallel} \equiv \frac{1}{\boldsymbol{q}_{\parallel}^{2}} \frac{d}{dt} (\boldsymbol{q}_{\parallel}^{2}) = \frac{e^{4} N_{i} L_{C}}{8 m_{p}^{2} c^{3} \boldsymbol{g}_{i}^{3} \boldsymbol{b}_{i}^{3} \boldsymbol{s}_{si} \boldsymbol{e}_{x}^{3/2} \boldsymbol{q}_{\parallel}^{2}} \sqrt{\frac{\boldsymbol{n}_{x}}{R}} \quad .$$

#### The IBS growth rate for transverse degree of freedom

- The heating of the longitudinal degree of freedom causes cooling for both transverse degrees of freedom;
- Additional mechanism heats the horizontal degree of freedom
  - ➤ At regions with non-zero dispersion, changes in longitudinal momentum change the particles reference orbits, which additionally excites the horizontal betatron motion,

$$\frac{d\mathbf{e}_{x}}{dt} = \left\langle A_{x} \frac{d\mathbf{q}_{\parallel}^{2}}{dt} \right\rangle_{s} ,$$
where
$$A_{x} = \frac{D^{2} + (D'\mathbf{b}_{x} + \mathbf{a}_{x}D)^{2}}{\mathbf{b}_{x}}$$

Finally, one can write for the emittance growth rates

$$\begin{bmatrix} \Lambda_{x} \\ \Lambda_{y} \end{bmatrix} \equiv \frac{1}{\boldsymbol{e}_{x,y}} \frac{d\boldsymbol{e}_{x,y}}{dt} = \frac{e^{4}N_{i}}{8\sqrt{2}m_{p}^{2}c^{3}\boldsymbol{g}_{i}^{3}\boldsymbol{b}_{i}^{3}\boldsymbol{s}_{si}\boldsymbol{e}_{x,y}}$$

$$\sqrt{\frac{1}{\sqrt{\boldsymbol{q}_{x}^{2} + \boldsymbol{q}_{y}^{2}}} \frac{L_{C}}{\boldsymbol{s}_{x}\boldsymbol{s}_{y}}} \begin{bmatrix} 2A_{x}\boldsymbol{\Xi}_{\parallel}(\boldsymbol{q}_{x},\boldsymbol{q}_{y}) - \frac{\boldsymbol{b}_{x}}{\boldsymbol{g}_{i}^{2}} \boldsymbol{\Xi}_{\perp}(\boldsymbol{q}_{x},\boldsymbol{q}_{y}) \\ -\frac{\boldsymbol{b}_{y}}{\boldsymbol{g}_{i}^{2}} \boldsymbol{\Xi}_{\perp}(\boldsymbol{q}_{y},\boldsymbol{q}_{x}) \end{bmatrix}$$
s

where

$$\Xi_{\perp}(x,y) \approx 1 + \frac{2\sqrt{2}}{\mathbf{p}} \ln \left( \frac{\sqrt{3x^2 + y^2}}{2y^2} x \right) + \frac{0.5429 \ln(y/x)}{\sqrt{1 + \ln^2(y/x)}}$$
.

- In the smooth focusing approximation,
- for equal horizontal and vertical emittances,  $e_x = e_y$ ,
- equal betatron tunes,  $\mathbf{n}_x = \mathbf{n}_y$ , and
- small contribution of energy spread into the beam size,  $D(\Delta p/p) \ll \mathbf{s}_x$ ,
  - ⇒ The emittance growth rate can be written in the following form

$$\Lambda_{x} = \frac{e^{4}N_{i}L_{c}}{16m_{p}^{2}c^{3}\boldsymbol{g}_{i}^{3}\boldsymbol{b}_{i}^{3}\boldsymbol{s}_{si}\boldsymbol{e}_{x}^{5/2}}\sqrt{\boldsymbol{n}_{x}}\left(\frac{2}{\boldsymbol{n}_{x}^{2}}-\frac{1}{\boldsymbol{g}_{i}^{2}}\right)$$

### 2. Intrabeam Scattering Estimate for the First Stage

#### Intra-beam scattering at injection

$$\theta_{\text{trinj}} = 3.177 \times 10^{-6}$$

$$\theta_{trinj} \cdot \beta_x = 0.055$$
 cm

$$\sigma_{sinj} = 15.221$$
 cm

$$\Delta PoverP_{inj} = 1.555 \times 10^{-3}$$

#### Intra-beam scattering at full energy

$$\theta_{\rm tr} = 6.738 \times 10^{-7}$$

$$\theta_{tr} \cdot \beta_{x} = 0.012$$
 cm

$$\sigma_{\rm S} = 6.747$$
 cm

$$\Delta$$
PoverP = 1.5× 10<sup>-4</sup>

$$\frac{\theta_{\text{trinj}}}{\Delta \text{PoverP}_{\text{inj}}} \cdot \gamma_{\text{inj}} = 1.959$$

$$L_c(\gamma_{inj}, \sigma_{sinj}) = 25.437$$

$$\frac{T_{longit}(\gamma_{inj}, \sigma_{sinj}, \Delta PoverP_{inj})}{360024} = 64.396 \text{ days}$$

$$\frac{T_{tr}(\gamma_{inj}, \sigma_{sinj})}{360024} = 6.231 \text{ days}$$

$$\frac{\theta_{tr}}{\Delta PoverP} \cdot \gamma = 95.758$$

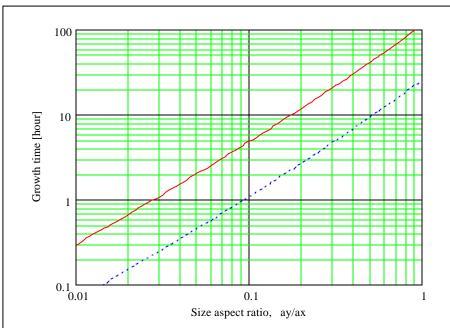
$$L_c(\gamma, \sigma_s) = 29.682$$

$$\frac{T_{longit}(\gamma, \sigma_s, \Delta Pover P)}{3600 24} = 23.846 \text{ days}$$

$$\frac{T_{tr}(\gamma, \sigma_s)}{360024} = 11.019 \text{ days}$$

## 3. Intrabeam Scattering Estimate for the Second Stage

- SR damping decreases beam sizes to the level where it will be stabilized by IBS or other heating mechanism
  ➤ SR damping time is ~ 2 hour
- We need to prevent decrease of horizontal and longitudinal emittances to avoid stronger IBS and TMCI.
  - Vertical beam size decrease is limited by increase in IBS for horizontal degree of freedom. That happens at the size aspect ratio,  $\mathbf{s}_y / \mathbf{s}_x$ , of about 0.15.
    - Vertical emittance needs to be kept at this level to prevent growth of horizontal emittance.



Growth rates for the longitudinal momentum spread (red solid line) and horizontal beam size (blue dashed line) as function of beam size aspect ratio,  $\mathbf{s}_y/\mathbf{s}_x$ ;  $\mathbf{s}_{\mathbf{D}p/p} = 5 \cdot 10^{-5}$ ,  $\mathbf{s}_s = 2.6$  cm,  $\mathbf{e}_{nx} = 0.2$  mm mrad

4. Conclusions

- IBS does not contribute into emittance growth for the first stage
- IBS is one of limiting factors for the second stage after the SR damping will cause significant cooling of beam emittances